

# Homework Set 5

Due: November 18, 2008, *before class*

## 1. White Dwarf Cooling

- (a) Following the lecture notes, compute the time scale on which the *effective* surface temperature changes.

Assume that radius of the white dwarf does not change.

$$L = 4\pi\sigma R^2 T_{\text{eff}}^4$$

⇒

$$\frac{d \ln L}{d \ln T_{\text{eff}}} = 4$$

Using the definitions

$$\tau_{\text{cool}} = -\frac{dt}{d \ln L} = \frac{3}{7} \frac{\mathcal{R}}{\mu_1 C_{\text{WD}}^{2/7}} \left(\frac{M}{L}\right)^{5/7} \approx 2.5 \times 10^6 \left(\frac{M}{M_{\odot}}\right)^{5/7} \left(\frac{L}{L_{\odot}}\right)^{-5/7} \text{ yr}$$

and analogously

$$\tau_{T_{\text{eff}}} = -\frac{dt}{d \ln T_{\text{eff}}}$$

we obtain

$$\begin{aligned} \tau_{T_{\text{eff}}} &= -\frac{dt}{d \ln T_{\text{eff}}} = -\frac{d \ln L}{d \ln T_{\text{eff}}} \frac{dt}{d \ln L} = 4 \tau_{\text{cool}} \\ \tau_{T_{\text{eff}}} &= \frac{12}{7} \frac{\mathcal{R}}{\mu_1 C_{\text{WD}}^{2/7}} \left(\frac{M}{L}\right)^{5/7} \approx 10^7 \left(\frac{M}{M_{\odot}}\right)^{5/7} \left(\frac{L}{L_{\odot}}\right)^{-5/7} \text{ yr} \end{aligned}$$

Score: 2

- (b) Compute the time scale on which the core temperature changes.

Similarly, using

$$L = M C_{\text{WD}} T_c^{7/2}$$

we have

$$\frac{d \ln L}{d \ln T_c} = 7/2$$

and therefore

$$\begin{aligned} \tau_{T_c} &= -\frac{dt}{d \ln T_c} = -\frac{d \ln L}{d \ln T_c} \frac{dt}{d \ln L} = \frac{7}{2} \tau_{\text{cool}} \\ \tau_{T_c} &= \frac{3}{2} \frac{\mathcal{R}}{\mu_1 C_{\text{WD}}^{2/7}} \left(\frac{M}{L}\right)^{5/7} \approx 9 \times 10^6 \left(\frac{M}{M_{\odot}}\right)^{5/7} \left(\frac{L}{L_{\odot}}\right)^{-5/7} \text{ yr} \end{aligned}$$

Score: 2

- (c) For the case of transition to a solid, compute the cooling time scale on which the luminosity changes for the case that specific heat is increased from  $\frac{3}{2} k_B/\text{ion}$  to  $3 k_B/\text{ion}$ . What is the signature in the WD luminosity function of such a change?

The specific heat capacity is doubled and therefore the cooling time scale becomes twice as long. The result is a bump in the WD luminosity function.

Score: 2

- (d) For the case of a cold solid, assume that the specific heat capacity per ion is  $DT^3$  where  $D$  is a constant.

**Compute the cooling time scale on which the luminosity changes for this case.**

For the internal energy we then have

$$U_1 = \frac{D}{\mu_1} MT_c^3$$

Following the lecture notes step by step we get

- luminosity equals loss in internal energy

$$L = -\frac{dU_1}{dt} = -\frac{D}{\mu_1} M \frac{dT_c^3}{dt} = -3 \frac{D}{\mu_1} MT_c^2 \frac{dT_c}{dt}$$

- using  $L \propto T_c^{7/2}$  (recall  $L = M C_{\text{WD}} T_c^{7/2}$ ) we can write

$$L = -\frac{6}{7} \frac{D}{\mu_1} M \frac{T_c^3}{L} \frac{dL}{dt}, \quad \frac{dL}{dt} = -\frac{7}{6} \frac{\mu_1}{D} \frac{L^2}{T_c^3 M}$$

- As before, define cooling time as time scale of luminosity drop by  $e$ :

$$\tau_{\text{cool}} = -\frac{dt}{d \ln L} = -L \left( \frac{dL}{dt} \right)^{-1} = \frac{6}{7} \frac{D}{\mu_1} \frac{M}{L} T_c^3$$

- eliminating  $T_c$  using  $L/M = C_{\text{WD}} T_c^{7/2}$  we obtain

$$\tau_{\text{cool}} = \frac{6}{7} \frac{D}{\mu_1 C_{\text{WD}}^{6/7}} \left( \frac{M}{L} \right)^{1/7} = \frac{6}{7} \frac{D}{\mu_1 C_{\text{WD}} T_c^{1/2}}$$

Without knowledge of  $D$  we cannot go any further. Note that the cooling rate now only very weakly depends on  $L$ ; it is almost constant  $\Rightarrow$  Fast cooling! Score: 6

## 2. Neutron Stars

Assume a non-rotating neutron star of gravitational mass  $1.4 M_\odot$  and radius 10 km.

The relativistic gravitational redshift  $(1+z)$  is given by

$$1+z = 1 / \sqrt{1 - R_s/r}$$

where

$$R_s = \frac{2GM}{c^2}$$

is the Schwarzschild radius. The relativistic gravitational acceleration is then given by

$$g_{\text{rel}} = (1+z) \frac{GM}{r^2} = \frac{GM}{r^2 \sqrt{1 - 2GM/(c^2 r)}}$$

- (a) **Compute, in classical approximation, the gravitational potential at the surface.**

$$\begin{aligned} \phi &= \frac{GM}{r} = \frac{G 1.4 M_\odot}{10^6 \text{ cm}} = \frac{6.67 \times 10^8 \text{ cm s}^{-2} \text{ g}^{-1} \times 1.4 \times 1.9891 \times 10^{33} \text{ g}}{10^6 \text{ cm}} \\ \phi &= 1.858 \times 10^{20} \text{ cm}^2 \text{ s}^{-2} = 0.2067 c^2 \end{aligned}$$

Score: 2

- (b) **Compute the relativistic gravitational potential at the surface of the neutron star.**

If something is redshifted by a factor  $1 + z$  when observed at infinity (large distance), this means that the original energy was higher by a factor  $1 + z$ ; what we observe is only  $1/(1 + z)$  of the original energy; the remainder,  $1 - 1/(1 + z) = z/(1 + z)$  has been lost. Due to energy conservation, this energy lost has to equal the potential energy, i.e.,  $\phi = -c^2 z/(1 + z)$ .

$$R_s = \frac{2GM}{r^2} = \frac{2 \times 6.67 \times 10^8 \text{ cm s}^{-2} \text{ g}^{-1} \times 1.4 \times 1.989 \times 10^{33} \text{ g}}{2.998 \times 10^{10} \text{ cm s}^{-1}} = 4.135 \times 10^5 \text{ cm}$$

$$\frac{R_s}{r} = \frac{4.135 \times 10^5 \text{ cm}}{10^6 \text{ cm}} = 0.4135$$

$$1 + z = \frac{1}{\sqrt{1 - R_s/r}} = \frac{1}{\sqrt{1 - 0.4135}} = 1.3058, \quad z = 0.3058$$

$$\phi = \frac{-c^2 z}{z + 1} = -c^2 \left(1 - \frac{1}{z + 1}\right) = -c^2 \left(1 - \frac{1}{1.3058}\right) = -c^2 \times 0.2342$$

$$\phi = -(2.998 \times 10^{10} \text{ cm s}^{-1})^2 \times 0.2342 = -8.988 \times 10^{20} \text{ cm}^2 \text{ s}^{-2} \times 0.2342$$

$$\phi = -2.105 \times 10^{20} \text{ cm}^2 \text{ s}^{-2}$$

Score: 4

- (c) **Show that for  $r \gg R_s$  the usual Newtonian formula for gravitational potential and acceleration is recovered.**

For  $x \ll 1$  use the approximation formulae

$$\sqrt{1 \pm x} \approx 1 \pm \frac{1}{2}x, \quad \frac{1}{x} \approx 1 - x$$

We note that for  $r \gg R_s$  we have  $R_s/r \ll 1$ . We then have

$$\phi = \frac{-c^2 z}{z + 1} = -c^2 \left(1 - \frac{1}{z + 1}\right) = -c^2 \left(1 - \sqrt{1 - \frac{R_s}{r}}\right)$$

$$\phi \approx -c^2 \left(1 - \left(1 - \frac{1}{2} \frac{R_s}{r}\right)\right) = -\frac{R_s c^2}{2r} = -\frac{2GMc^2}{2rc^2} = -\frac{GM}{r}$$

And for acceleration

$$1 + z = \frac{1}{\sqrt{1 - \frac{R_s}{r}}} \approx 1 + \frac{1}{2} \frac{R_s}{r} \rightarrow 1$$

$\Rightarrow$

$$g_{\text{rel}} = (1 + z) \frac{GM}{r^2} \approx \frac{GM}{r^2}$$

Score: 4

- (d) Assume the neutron star accretes material composed of solar composition. For simplicity, assume this is 75 % hydrogen and 25 % helium, by mass fraction.

**Compare the specific energy that can be released by nuclear burning to  $^{56}\text{Fe}$  to the gravitational potential at the surface of the neutron star.**

For simplicity we neglect the energy lost due to neutrinos.

mass excess of

- $^1\text{H}$ : 7.289 MeV
- $^4\text{He}$ : 2.435 MeV
- $^{56}\text{Fe}$ : -60.601 MeV

mass excess per nucleon:

- ${}^1\text{H}$ : 7.289 MeV
- ${}^4\text{He}$ :  $2.435 \text{ MeV}/4 = 0.60875 \text{ MeV}$
- ${}^{56}\text{Fe}$ :  $-60.601 \text{ MeV}/56 = -1.08216 \text{ MeV}$

total energy release per nucleon from

- ${}^1\text{H}$ :  $e_{1\text{H}} = 7.289 \text{ MeV} - (-1.08216 \text{ MeV}) = 8.3712 \text{ MeV}$
- ${}^4\text{He}$ :  $e_{4\text{He}} = 0.60875 \text{ MeV} - (-1.08216 \text{ MeV}) = 1.6909 \text{ MeV}$

total energy per nucleon release weighed by mass fractions  $X = 0.75$  and  $Y = 0.25$

$$e = X \times e_{1\text{H}} + Y \times e_{4\text{He}} = 0.75 \times 8.3712 \text{ MeV} + 0.25 \times 1.6909 \text{ MeV} = 6.7011 \text{ MeV}$$

energy release per gram is the approximately

$$e = 6.7011 \text{ MeV} \times \left( \frac{1.6021773 \times 10^{-6} \text{ erg}}{1 \text{ MeV}} \right) \times N_A = 1.0736 \times 10^{-5} \text{ erg} \times 6.0221367 \times 10^{23} \text{ g}^{-1}$$

$$e = 6.4656 \times 10^{18} \text{ erg g}^{-1} = 7.1939 \times 10^{-3} c^2$$

**For GR gravity:**

The nuclear energy is only  $e/\phi = 7.1939 \times 10^{-3} c^2 / 0.2342 c^2 \approx 3.07\%$  of the gravitational potential energy.

**For Newtonian gravity:**

The nuclear energy is only  $e/\phi = 7.1939 \times 10^{-3} c^2 / 0.2067 c^2 \approx 3.48\%$  of the gravitational potential energy.

Score: 4

- (e) **If the energy released by nuclear burning was used to lift off the surface of the neutron star, how far can you lift it, or how fast can the material still move at infinity (depending on which of the energies of the previous questions is greater)?**

In principles one needs radius coordinate  $r$  as a function of potential  $\phi$ , i.e.,  $r(\phi)$ , and can then compute

$$dr = r(\phi + e) - r(\phi)$$

. There are two possibilities:

- i. Since  $e \ll |\phi|$  we could use

$$dr \approx -e \left( \frac{dr}{d\phi} \right) = -e \left( \frac{d\phi}{dr} \right)^{-1}$$

**For Newtonian gravity:**

$$\frac{d\phi}{dr} = \frac{d}{dr} \left( -\frac{GM}{r} \right) = \frac{GM}{r^2}$$

and we obtain

$$dr \approx \frac{er^2}{GM} = \frac{2er^2}{c^2 R_s} = \frac{2 \times 7.1939 \times 10^{-3} \times (10^6 \text{ cm})^2}{4.135 \times 10^5 \text{ cm}} = 34796 \text{ cm}$$

$$dr \approx 350 \text{ m}$$

**For GR gravity:**

$$\frac{d\phi}{dr} = -c^2 \frac{d}{dr} \left( 1 - \sqrt{1 - R_s/r} \right) = \frac{c^2 R_s}{2r^2} \left( 1 - \frac{R_s}{r} \right)^{-1/2}$$

and we obtain

$$dr \approx \frac{2er^2}{c^2 R_s} \left( 1 - \frac{R_s}{r} \right)^{1/2} = \frac{2 \times 7.1939 \times 10^{-3} \times (10^6 \text{ cm})^2}{4.135 \times 10^5 \text{ cm}} (1 - 0.4135)^{1/2} = 26647 \text{ cm}$$

$$dr \approx 270 \text{ m}$$

ii. Exact solution.

For Newtonian gravity:

$$e = GM \left( \frac{1}{r} - \frac{1}{r + dr} \right)$$

$$dr = r \left( \frac{1}{1 - er/GM} - 1 \right) = r \left( \frac{1}{1 - 2er/R_s c^2} - 1 \right)$$

$$\frac{2er}{R_s c^2} = \frac{2 \times 7.1939 \times 10^{-3} \times 10^6 \text{ cm}}{4.135 \times 10^5 \text{ cm}} = 3.4796 \times 10^{-2}$$

$$dr = 10^6 \text{ cm} \times \left( \frac{1}{1 - 3.4796 \times 10^{-2}} - 1 \right) = 36050 \text{ cm} \approx 360 \text{ m}$$

For GR gravity:

$$e = -c^2 \left( 1 - \sqrt{1 - R_s/r} - \left( 1 - \sqrt{1 - R_s/(r + dr)} \right) \right)$$

$$e = c^2 \left( \sqrt{1 - R_s/(r + dr)} - \sqrt{1 - R_s/r} \right)$$

$$dr = R_s \left[ 1 - \left( \sqrt{1 - R_s/r} + \frac{e}{c^2} \right)^2 \right]^{-1} - r = R_s \left[ 1 - \left( \frac{1}{1+z} + \frac{e}{c^2} \right)^2 \right]^{-1} - r$$

$$dr = 413492 \text{ cm} \left[ 1 - \left( \frac{1}{1.3057} + 7.1939 \times 10^{-3} \right)^2 \right]^{-1} - 10^6 \text{ cm} = 27647 \text{ cm} \approx 280 \text{ m}$$

Score: 4

### 3. Black Holes

(a) Compute the Schwarzschild Radius of the sun and of the earth.

$$R_s = \frac{2GM}{c^2} = M \times \frac{2G}{c^2} = M \times \frac{2 \times 6.67 \times 10^{-8} \text{ cm s}^2 \text{ g}^{-1}}{(2.9979 \times 10^{10} \text{ cm s}^{-1})^2} = M \times 1.4842 \times 10^{-28} \text{ cm g}^{-1}$$

$$R_{s,\odot} = M_{\odot} \times \frac{2G}{c^2} = 1.9891 \times 10^{33} \text{ g} \times 1.4842 \times 10^{-28} \text{ cm g}^{-1} = 2.95 \times 10^5 \text{ cm} \approx 3 \text{ km}$$

$$R_s = 2.954 \times 10^5 \text{ cm} \left( \frac{M}{M_{\odot}} \right)$$

$$R_{s,\oplus} = M_{\oplus} \times \frac{2G}{c^2} = 5.9742 \times 10^{27} \text{ g} \times 1.4842 \times 10^{-28} \text{ cm g}^{-1} = 0.8867 \text{ cm}$$

Score: 4

(b) In classical approximation using Keplerian mechanics, at what orbital distance would an object be from the center of the Sun assumed as a point mass that has an orbital velocity of the speed of light?

In Keplerian orbit centrifugal force equals gravity:

$$\frac{GM_{\odot}}{r^2} = \frac{v^2}{r}$$

for  $v = c$  one obtains for  $r$ :

$$r = \frac{GM_{\odot}}{c^2} = \frac{1}{2} R_{s,\odot} = 1.5 \times 10^5 \text{ cm} = 1.5 \text{ km}$$

Score: 4

**Please use cgs units for calculations and numerical values.**